CALCULATING THE POSITRON – ELECTRON CORRELATION ENERGY IN ZnO WITH THE MODIFIED SINGLE WAVE FUNCTION FOR POSITRON

Chau Van Tao⁽¹⁾, Trinh Hoa Lang⁽¹⁾, Nguyen Anh Tuan⁽²⁾, Le Hoang Chien⁽¹⁾, Nguyen Huu Loc⁽¹⁾

(1) University of Science

(2) Research and Development Center for Radiation Technology

ABSTRACT: The ZnO – positron system is studied and its positron – electron correlation energy is estimated in its ground state. The positron binds with the outer shell electrons of Zinc and Oxigen to form the pseudo ZnO – positron molecule before it anihilates with one of these electrons. In this work, the single wave function for positron is modified according to the principle of linear superposition, and by using Variational Quantum Monte – Carlo method (VQMC) the correlation energy of this system is estimated with the value $E_c^{e-p} = -9.3 \pm 1.1$ eV. It turns out that the value is closer to results estimated by other methods than the value that we had done before.

Keywords: Positron, Variational Quantum Monte - Carlo

INTRODUCTION

There are many methods that have been used to estimate the positron-electron correlation energy [2]. In the article, we use the VQMC method to estimate the positronelectron correlation energy. The method is more successful when we make the form of the trial wave functions of particles in a considered system are loser to that of the exact wave function. So, our main aim is to modify the form of the trial wave function of the positron.

THEORY

Variational Quantum Monte Carlo method

The quantum mechanical system is represented by an exact wave function $\psi(\mathbf{R})$ and the average value of the system's energy (we want to get the quantity in the process of estimating the electron, positron correlation – energy of the zinc oxide - positron system) is given by

$$\left\langle \mathbf{H}\right\rangle = \frac{\int d\mathbf{R}\psi^{*}(\mathbf{R})\mathbf{H}(\mathbf{R})\psi(\mathbf{R})}{\int d\mathbf{R}\psi^{*}(\mathbf{R})\psi(\mathbf{R})}$$

However, we cannot properly construct the form of the exact wave function. This also means that we cannot exactly calculate H(R) in theory. Therefore, we must choose a trial total wave function $\psi_T(R, \alpha)$ that depends on a set of parameters α . In the circumstance, we define two new quantities

$$\rho(\mathbf{R}, \alpha) = \frac{\left|\psi_{\mathrm{T}}(\mathbf{R}, \alpha)\right|^{2}}{\int d\mathbf{R} \left|\psi_{\mathrm{T}}(\mathbf{R}, \alpha)\right|^{2}} = \left|\psi_{\mathrm{T}}(\mathbf{R}, \alpha)\right|^{2} \qquad (2)$$

$$E(\mathbf{R}, \alpha) = \frac{H(\mathbf{R})\psi_{T}(\mathbf{R}, \alpha)}{\psi_{T}(\mathbf{R}, \alpha)} \qquad (3)$$

With combining (1), (2) and (3) we can figure out the equation (4)

$$\left\langle E \right\rangle = \frac{\int dR \rho(R, \alpha) \frac{H(R)\psi_{T}(R, \alpha)}{\psi_{T}(R, \alpha)}}{\int dR\psi_{T}^{*}(R, \alpha)\psi_{T}(R, \alpha)} \qquad (4)$$

can approach $\langle H \rangle$ only if the trial wave function is close to the exact wave function. To satisfy this condition, we can vary values of α and search for the minimum value of $\langle E \rangle$ which corresponds to the exact wave function. The calculations are solved by using techniques from the Variational Quantum Monte Carlo method.

Model of the zinc oxide - positron

In this paper, the electron, positron correlation – energy of the zinc oxide - positron system which shown in figure 1 is estimated by applying to the VMC method in which the Hamiltonian and the trial total wave function were used.

With the Born – Oppenheimer approximation [3], the Hamiltonian for the system is then

$$H = -\frac{1}{2} \sum_{i=1}^{6} \nabla_{i}^{2} - \frac{1}{2} \nabla_{p}^{2} + \frac{1}{2} \sum_{i=1}^{6} \sum_{\substack{j=1\\j\neq i}}^{6} \frac{1}{|\mathbf{r}_{i} - \mathbf{r}_{j}|} + \frac{1}{2} \frac{Z_{Zn} Z_{0}}{|\mathbf{d}_{Zn} - \mathbf{d}_{0}|} + \frac{1}{2} \frac{Z_{Zn} Z_{0}}{|\mathbf{d}_{Zn} - \mathbf{d}_{0}|} - \sum_{i=1}^{6} \frac{Z_{Zn}}{|\mathbf{r}_{i} - \mathbf{r}_{i}|} - \sum_{i=1}^{6} \frac{Z_{Zn}}{|\mathbf{r}_{i} - \mathbf{r}_{i}|} + \frac{Z_{Zn}}{|\mathbf{r}_{p} - \mathbf{r}_{i}|} + \frac{Z_{Op}}{|\mathbf{r}_{p0}|} - \sum_{i=1}^{6} \frac{1}{|\mathbf{r}_{ip}|}$$

$$(5)$$

where

- $-\frac{1}{2}\nabla_i^2$, $-\frac{1}{2}\nabla_p^2$ is alternatively the

kinetic operators of the i-th electron and the positron.

- r_{iZn} (r_{iO}) is the distance from the position of the i-th electron to the position of zinc atom (oxygen atom).
- r_{pZn} (r_{pO}) is the distance from the position of the positron to the position of zinc atom (oxygen atom). Similarly, r_{ip} is the

distance from the position of the i-th electron to the position of the positron.

- r_i (r_j) is the distance from the position of the i-th (j-th) electron to the coordinate angle.
- d_{Zn} (d_O) is the distance from the position of zinc atom (oxygen atom) to the coordinate angle.
- Z_{Zne} (Z_{Oe}) is the effective nuclear charge of zinc (oxygen) atom for the electrons.
- Z_{Znp} (Z_{Op}) is the effective nuclear charge of zinc (oxygen) atom for the positron.

 $Z_{\rm O}$ is the effective nuclear charge of oxygen atom for zinc atom. Conversely, $Z_{\rm Zn}$ is

the nuclear charge of zinc atom for oxygen atom.



Fig.1. The schematic of the zinc oxide - positron system

The trial total wave function for the system used in this work is a product of single – particle wave functions multiplied by a Jastrow factor (an exponent of many particle correlation factors)

$$\Psi_{T} = \prod_{i=1}^{4} \Psi_{io} \prod_{i=1}^{2} \Psi_{iZn} \prod_{\substack{i=1\\j>i}}^{6} \Psi^{e \cdot e}_{ij} \prod_{i=1}^{6} \Psi^{e \cdot p}_{ip} \Psi_{p}$$
(6)

with

 Ψ_{iO} is the single – particle wave function for the i-th electron belonging to oxygen atom. With the Slater's approach [4], it takes the following form

-
$$\psi_{i0} = N_0 r_{i0} e^{-\lambda_0 r_{i0}} + N_0 r_{iZn} e^{-\lambda_{Zn} r_{iZn}}$$
 (7)

- Ψ_{iZn} is the single – particle wave function for the i-th electron belonging to zinc atom. With the Slater's approach, it takes the following form

-
$$\psi_{iZn} = N_{Zn} r^4_{iZn} e^{-\lambda_{Zn} r_{iZn}} + N_{Zn} r^4_{io} e^{-\lambda_o r_{io}} (8)$$

Ψ^{e-e}_{ij} is the Jastrow factor that reflects the correlation between the i-th electron and the j-th electron. With the Pade's approach [1], it takes the form

$$\psi^{e-e}_{ij} = e^{\frac{\beta r_{ij}}{1 + \alpha r_{ij}}} \quad (9)$$

- Ψ_{ip}^{e-p} is the Jastrow factor that reflects the correlation between the i-th electron and the positron. With the Pade's approach, it takes the form

$$- \qquad \psi^{e-p}_{ip} = e^{\frac{\beta r_{ip}}{1 + \alpha r_{ip}}}$$
(10)

 ψ_p is the single – particle wave function for the positron. When the positron moves into Zinc oxide molecule, the positron may belong to either zinc atom or oxygen atom. So, we assume that positron exists in some allowed state supported by either the nucleus of Zinc atom or the nucleus of oxygen atom. According to the principle Trang 79 of linear superposition, the single – particle wave function for positron takes form

Applying the Hamiltonian in eq. (5) to eq.(6), the total energy takes the following form:

$$\Psi_{p} = \frac{1}{2} N_{p} r_{p0} e^{-\lambda_{p0} r_{p0}} + \frac{1}{2} N_{p} r_{pZn}^{4} e^{-\lambda_{pZn} r_{pZn}}$$
(11)
$$E = \sum_{i=1}^{6} (2K_{i} - F_{i}^{2}) + 2K_{p} - F_{p}^{2} + V$$
(12)

where

$$K_{i} = -\frac{1}{4} \nabla_{i}^{2} \ln \psi_{T} = -\frac{1}{4} \left[\nabla_{i}^{2} \left(\ln \varphi_{e} \right) + \nabla_{i}^{2} \left(\sum_{\substack{j=1\\j\neq i}}^{6} \ln \psi_{ij}^{e_e} \right) + \nabla_{i}^{2} \left(\ln \psi_{ip}^{e_p} \right) \right]$$
(13)
$$F_{i} = \frac{1}{\sqrt{2}} \overrightarrow{\nabla_{i}} \ln \psi_{T} = \frac{1}{\sqrt{2}} \left(\overrightarrow{\nabla_{i}} \ln \varphi_{e} + \overrightarrow{\nabla_{i}} \sum_{\substack{j=1\\j\neq i}}^{6} \ln \left(\psi_{ij}^{e_e} \right) + \overrightarrow{\nabla_{i}} \ln \left(\psi_{ip}^{e_p} \right) \right)$$
(14)

$$K_{p} = -\frac{1}{4}\nabla_{p}^{2}\ln\psi_{T} = -\frac{1}{4}\left[\nabla_{p}^{2}\left(\ln\psi_{p}\right) + \nabla_{p}^{2}\left(\sum_{i=1}^{6}\ln\psi_{ip}^{e_{p}}\right)\right]$$
(15)

$$F_{p} = \frac{1}{\sqrt{2}} \overrightarrow{\nabla_{p}} ln \psi_{T} = \frac{1}{\sqrt{2}} \left(\overrightarrow{\nabla_{p}} \left(\sum_{i=1}^{6} ln \left(\psi_{ip}^{e_{p}} \right) \right) + \overrightarrow{\nabla_{p}} ln \left(\psi_{p} \right) \right)$$
(16)

$$V = \frac{Z_{Zn}Z_{O}}{\left|d_{Zn} - d_{O}\right|} + \frac{1}{2}\sum_{i=1}^{6}\sum_{\substack{j=1\\j\neq i}}^{6}\frac{1}{\left|r_{i} - r_{j}\right|} - \sum_{i=5}^{6}\frac{Z_{Zne}}{\left|r_{iZn}\right|} - \sum_{i=1}^{4}\frac{Z_{Oe}}{\left|r_{iO}\right|} + \frac{Z_{Znp}}{\left|r_{pZn}\right|} + \frac{Z_{Op}}{\left|r_{pO}\right|} - \sum_{i=1}^{6}\frac{1}{\left|r_{ip}\right|}$$
(17)

The form of the electron, positron correlation energy which is extract from (12) is as then

$$E^{e_{p}} = \sum_{i=1}^{6} \left(2K_{i}^{e_{p}} - \left(F_{i}^{e_{p}}\right)^{2} \right) + 2K_{p}^{e_{p}} - \left(F_{p}^{e_{p}}\right)^{2}$$
(18)

where

$$K_{i}^{e_{p}} = -\frac{1}{4} \nabla_{i}^{2} \left(ln \psi_{ip}^{e_{p}} \right) \text{ and } K_{p}^{e_{p}} = -\frac{1}{4} \nabla_{p}^{2} \left(\sum_{i=1}^{6} ln \psi_{ip}^{e_{p}} \right)$$
(19)

Trang 80

$$\left(F_{i}^{e_p}\right)^{2} = \frac{1}{2} \left(2\overrightarrow{\nabla_{i}}\ln\left(\psi_{ip}^{e_p}\right)\overrightarrow{\nabla_{i}}\ln\varphi_{e} + 2\overrightarrow{\nabla_{i}}\ln\left(\psi_{ip}^{e_p}\right)\overrightarrow{\nabla_{i}}\sum_{\substack{j=1\\j\neq i}}^{6}\ln\left(\psi_{ij}^{e_e}\right) + \left[\overrightarrow{\nabla_{i}}\ln\left(\psi_{ip}^{e_p}\right)\right]^{2}\right)$$
(20)
$$\left(F_{p}^{e_p}\right)^{2} = \frac{1}{2} \left(2\overrightarrow{\nabla_{p}}\left(\sum_{i=1}^{6}\ln\left(\psi_{ip}^{e_p}\right)\right)\overrightarrow{\nabla_{p}}\ln\left(\psi_{p}\right) + \left[\overrightarrow{\nabla_{p}}\left(\sum_{i=1}^{6}\ln\left(\psi_{ip}^{e_p}\right)\right)\right]^{2}\right)$$
(21)
$$V^{e_p} = -\sum_{i=1}^{6}\frac{1}{\left|r_{ip}\right|}$$
(22)

CALCULATING

In calculation, we will generate 8 sets of data by varying one of the eight parameters λ_{Zn} , λ_O , α , β , λ_{pZn} , λ_{pO} , α' , β' and keeping the other seven parameters as constants. Runs were performed with N = 300 walkers and MCSteps = 10000 times (Monte Carlo steps).



Fig.2a. The average energy is in the term of λ_{Zn}



Fig.2b. The relative error is in the term of λ_{Zn}

Varying the parameter λ_{Zn} : First, we will vary the parameter λ_{Zn} from $\lambda_{ZnMin} = 3$ to $\lambda_{ZnMax} = 6.8$ with $\delta\lambda_{Zn} = 0.2$ while $\lambda_O = 3.2$, $\alpha = 0.14$, $\beta = 0.5$, $\lambda_{pZn} = 0.21$, $\lambda_{pO} = 0.72$, $\alpha' = 0.88$ and $\beta' = -1.5$ are considered as constants. The following plots show results for the average energy $\langle E \rangle$ and its relative error as functions of the variational parameter λ_{Zn} .

The average energy and the relative error is minimum at $\lambda_{Zn} = 4.2$ which corresponds to the optimized solution for the ground state.

Varying the parameter λ_0 : Second, we will vary the parameter λ_0 from $\lambda_{OMin} = 4$ to $\lambda_{OMax} =$ 6.8 with $\delta\lambda_0 = 0.2$ while $\lambda_{Zn} = 4.2$, $\alpha = 0.14$, β = 0.5, $\lambda_{pZn} = 0.21$, $\lambda_{pO} = 0.72$, $\alpha' = 0.88$ and β' = -1.5 are considered as constants. The following plots show results for the average energy <E> and its relative error as functions of the variational parameter λ_0 .



Fig.3a. The average energy is in the term of λ_0



Fig.3b. The relative error is in the term of λ_0

The average energy and the relative error is minimum at $\lambda_0 = 5.2$ which corresponds to the optimized solution for the ground state.

Varying the parameter β : Next, we will vary the parameter β from $\beta_{Min} = 0.32$ to $\beta_{Max} =$ 0.7 with $\delta\beta = 0.02$ while $\lambda_O = 5.2$, $\lambda_{Zn} = 4.2$, α = 0.14, $\lambda_{pZn} = 0.21$, $\lambda_{pO} = 0.72$, $\alpha' = 0.88$ and β' = -1.5 are considered as constants. The following plots show results for the average energy <E> and its relative error as functions of the variational parameter β .



Fig.4a. The average energy is in the term of β





The average energy and the relative error is minimum at $\beta = 0.58$ which corresponds to the optimized solution for the ground state.

Varying the parameter α : Next, we will vary the parameter α from $\alpha_{Min} = 0.06$ to $\alpha_{Max} =$

Trang 82

0.48 with $\delta \alpha = 0.02$ while $\lambda_0 = 5.2$, $\lambda_{Zn} = 4.2$, $\beta = 0.58$, $\lambda_{pZn} = 0.21$, $\lambda_{pO} = 0.72$, $\alpha' = 0.88$ and $\beta' = -1.5$ are considered as constants. The following plots show results for the average energy <E> and its relative error as functions of the variational parameter α .



Fig.5a. The average energy is in the term of α



Fig.5b. The relative error is in the term of α

The average energy and the relative error is minimum at $\alpha = 0.18$ which corresponds to the optimized solution for the ground state.

Varying the parameter λ_{pZn} : Next, we will vary the parameter λ_{pZn} from $\lambda_{pZnMin} = 0.06$ to $\lambda_{pZnMax} = 0.48$ with $\delta\lambda_{pZn} = 0.02$ while $\lambda_{Zn} = 4.2$, $\lambda_0 = 5.2$, $\alpha = 0.18$, $\beta = 0.58$, $\lambda_{pO} = 0.72$, $\alpha' = 0.88$ and $\beta' = -1.5$ are considered as constants. The following plots show results for the average energy $\langle E \rangle$ and its relative error as functions of the variational parameter λ_{pZn} .



Fig.6a. The average energy is in the term of λ_{pZn}



Fig.6b. The relative error is in the term of λ_{Zn}

The average energy and the relative error is minimum at $\lambda_{pZn} = 0.22$ which corresponds to the optimized solution for the ground state.

Varying the parameter λ_{pO} : Next, we will vary the parameter λ_{pO} from $\lambda_{pOMin} = 0.1$ to $\lambda_{pOMax} = 0.48$ with $\delta\lambda_{pO} = 0.02$ while $\lambda_{Zn} = 4.2$, $\lambda_O = 5.2$, $\alpha = 0.18$, $\beta = 0.58$, $\lambda_{pZn} = 0.22$, $\alpha' = 0.88$ and $\beta' = -1.5$ are considered as constants. The following plots show results for the average energy <E> and its relative error as functions of the variational parameter λ_{pO} .



Fig.7a. The average energy is in the term of λ_{pO}



Fig.7b. The relative error is in the term of λ_{pO}

The average energy and the relative error is minimum at $\lambda_{pO} = 0.14$ which corresponds to the optimized solution for the ground state. Varying the parameter β' : Next, we will vary the parameter β' from $\beta'_{Min} = 0.02$ to β'_{Max} = 0.38 with $\delta\beta' = 0.02$ while $\lambda_{Zn} = 4.2$, $\lambda_O = 5.2$, $\alpha = 0.18$, $\beta = 0.58$, $\lambda_{pZn} = 0.22$, $\lambda_{pO} = 0.14$ and $\alpha' = 0.88$ are considered as constants. The following plots show results for the average energy <E> and the variance as functions of the variational parameter β' .

The average energy and the relative error is minimum at $\beta' = 0.04$ which corresponds to the optimized solution for the ground state.



Fig.8a. The average energy is in the term of β'



Fig.8b. The relative error is in the term of β'

Varying the parameter α : Next, we will vary the parameter α from $\alpha'_{Min} = 0.2$ to $\alpha'_{Max} =$ 0.38 with $\delta \alpha' = 0.2$ while $\lambda_{Zn} = 4.2$, $\lambda_O = 5.2$, α = 0.18, $\beta = 0.58$, $\lambda_{pZn} = 0.22$, $\lambda_{pO} = 0.14$ and β' = 0.04 are considered as constants. The following plots show results for the average energy <E> and its relative error as functions of the variational parameter α' .



Trang 83

Fig.9a. The average energy is in the term of α



Fig.9b. The relative error is in the term of α

The average energy and the relative error is minimum at $\alpha' = 1.2$ which corresponds to the optimized solution for the ground state.

Finally, we get the values of the optimized parameters listed in the following table 1

Table 1. The values of exact parameters

λο	λ_{Zn}	α	β	λ_{pZn}	λ_{pO}	a	β [`]
5.2	4.2	0.18	0.58	0.22	0.14	1.2	0.04

After varying alternately the parameters, we continuous to estimate the correlation – energy E^{e-p} with the set of the optimized parameters by applying to the formula (18) and the result is $E^{e-p} = -9.3 \pm 1.1$ (eV).

CONCLUSION

In this article, we described the Variational Quantum Monte Carlo method, the technique that was used to estimate the value of the electron, positron correlation - energy of ZnO molecule. The Hamiltonian and the many electron, positron wave function were also discussed.

With building a code based on programming language C++, performing the configuration with 300 walkers and 10000 MCSteps, the electron, positron correlation energy of ZnO molecule was estimated, $E^{e-p} = -9.3 \pm 1.1$ (eV). It turns out that the value is closer to results estimated by other methods [2] than the value that we had done before.

KHẢO SÁT NĂNG LƯỢNG TƯỜNG QUAN ELECTRON – POSITRON TRONG ZnO VỚI SỰ HIỆU CHỈNH HÀM SỐNG POSITRON

Châu Văn Tạo⁽¹⁾, Trịnh Hoa Lăng⁽¹⁾, Nguyễn Anh Tuấn⁽²⁾, Lê Hoàng Chiến⁽¹⁾, Nguyễn Hữu Lộc (1) Trường Đại học Khoa học Tự nhiên (2) Trung tâm Nghiên cứu và Triển khai Công nghệ Bức xạ

TÓM TÅT: Trong bài báo này, chúng tôi tính năng lượng tương quan electron – positron trong phân tử kẽm oxit (ZnO), trong đó giả thiết rằng positron liên kết với các electron thuộc phân lớp ngoài cùng của các nguyên tử kẽm và oxi trước khi nó hủy với một trong các electron đó. Với việc sử dụng phương pháp biến phân Monte - Carlo lượng tử (VQMC), đồng thời hiệu chỉnh hàm sóng của positron theo nguyên lý chồng chất nhiều trạng thái, năng lượng tương quan electron–positron được tính toán và nó có giá trị là $E_c^{e-p} = -9,3 \pm 1,1$ eV. Kết quả này (-9,3 eV) gần với các kết quả của một số chất được tính bởi các phương pháp khác hơn so với trường hợp mà chúng tôi đã tính trước đây.

Từ khóa: Positron, biến phân Monte - Carlo lượng tử

REFERENCES

- [1]. Louisa Màiri Fraser, Coulomb interactions and Positron annihilation in many fermion systems: A Monte Carlo approach, Thesis submitted for the degree of Doctor of Philosophy of the University of London and the Diploma of Imperial College, Technology and Medicine London (9/1995).
- [2]. M. J. Puska and R. M. Nieminen, Defect spectroscopy with positron: a general calculation method, *J. Phys. F: Met. Phys*, 333-346 (1983).
- [3]. N. W. Ashcroft and N. D. Mermin, Solid State Physics, (Holt-Saunders, 1976).
- [4]. Valerio Magnasco, Methods of Molecular Quantum Mechanics, University of Genoa, Genoa, Italy (2009).